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#### Four-State Modulation Continuous Variable Quantum Key Distribution over a 30-km Fiber and Analysis of Excess Noise \*

WANG Xu-Yang(王旭阳), BAI Zeng-Liang(白增亮), WANG Shao-Feng(王少锋), LI Yong-Min(李永民)\*\*, PENG Kun-Chi(彭堃墀)

State Key Laboratory of Quantum Optics and Quantum Optics Devices, Institute of Opto-Electronics, Shanxi University, Taiyuan 030006

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We report a fiber-based four-state discrete modulation continuous variable quantum key distribution system based on homodyne detection. A secret key rate of 1 kbit/s is achieved at a transmission distance of 30.2 km. Two factors that result in the excess noises of the quantum key distribution system are analyzed. The first is the relative phase dithering between the signal and local fields, and the second is the local field leakage into the signal field due to the scattering process that depolarizes the local field. It is found that the latter has a significant impact on the excess noise, which is the main limiting factor to the long-distance secure quantum transmission. Some protocols are also given to decrease the excess noise effectively.

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Coherent-state continuous variable quantum key distribution (CVQKD) protocols based on quadratures of the quantized electromagnetic field have seen significant progress during the last decade. Various kinds of protocols have been proposed, [1-17] some of which have been implemented experimentally and field testing was also performed recently.<sup>[18-25]</sup> Compared with its counterpart discrete variable quantum key distribution (DVQKD), CVQKD has some technological advantages such as better efficiency of homodyne detection over single photon counting at the telecom wavelength.

To achieve a long-distance CVQKD, two approaches have been proposed. One is based on a discrete modulation technique which includes two-state and four-state protocols, [10,13] etc. Such protocols allow high reconciliation efficiency for SNR close to 0. Meanwhile, the corresponding modulation procedure is simpler than that of Gaussian modulation and thus more practical. The other is to build a good reconciliation code with high efficiency for a Gaussian modulation coherent state protocol at a low SNR and it is shown that a long distance of 120 km can be realized for reasonable physical parameters.<sup>[16]</sup> Recently, a 24-km fiber-based discretely signaled CVQKD system using the technique of reverse reconciliation and postselection was reported, in which the security against the collective attack was guaranteed by quantum state tomography on a subset of Bob's data.<sup>[23]</sup> Experimental implementation of a no-switching discrete modulation protocol was also presented in free space.<sup>[22]</sup>

In this Letter, we present a fiber-based four state discrete modulation CVQKD system. It is found that the relative phase fluctuations between the signal and local fields, and local field leakage into the signal field due to the depolarized scattering process can cause a significant amount of excess noises, in particular

the latter one, which degrades the performance of the long-distance quantum distribution significantly.

The unconditional security of the four-state protocol based on homodyne detection has been proved and the protocol can be briefly described as follows.<sup>[10]</sup> Alice prepares and sends randomly one of the four coherent states:  $|\alpha \exp[i(2k+1)\pi/4]\rangle$  with  $k \in \{0, 1, 2, 3\}$ (here  $\alpha$  is taken as a real number). The receiver Bob measures randomly one of the quadratures  $\hat{X}$  or  $\hat{P}$  and obtains his results. The sign of Bob's results encodes the bit of the raw key.

The experimental setup of the fiber-based fourstate CVQKD is shown in Fig. 1.<sup>[19,21]</sup> The output beam of a continuous-wave 1550 nm single frequency fiber laser is modulated to 100-ns-wide pulses using the two cascaded high extinction ratio intensity modulators which are driven by a pulse generator at a repetition rate of 500 kHz. The optical pulses are divided into a strong local oscillator and a weak signal field by a 99/1 asymmetric fiber coupler. With variable attenuator and computer-driven intensity modulators and phase modulators, the coherent states can be prepared with random distribution in phase space as shown in Fig. 2. Through time multiplexing which is achieved by an 80 m fiber and polarization multiplexing realized by a PBS, the signal and local pulses are directed to a 30-km single mode fiber coil. After long-distance transmission the receiver Bob measures the quadrature amplitude or phase of the signal field randomly using a PHD. The phase modulator is used for randomly switching the measurement basis and locking the relative phase between the signal and local fields.

To facilitate the key transmission, the pulses are split into blocks. For the pulse repetition rate of  $500 \,\mathrm{kHz/s}$ , the time duration of each block is  $100 \,\mathrm{ms}$ , which is also the calibration time that we lock the relative phase between the local and signal fields. The

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relative phase for each block can be determined in real time by using a portion of the block pulses. In order to synchronize the pulses between Alice and Bob accurately, the whole system shares one common clock source generated by a data acquisition module (NI, USB 6259) located on Alice's site. On Bob's site, a 50/50 coupler is used to extract half the beam to recover the clock. To ensure the peak value of the PHD output that can be acquired accurately, a delay box is connected to the recovered clock signal to generate the desired time delay.



**Fig. 1.** The schematic diagram of the experimental setup. DAQ1, DAQ2, data acquisition module; AM, amplitude modulator; PM, phase modulator; SM, single mode; PHD, pulsed homodyne detector; PBS, polarizing beamsplitter.



**Fig. 2.** Conceptual schematic of the four-state protocol in phase space. Here  $\varepsilon_s$  is the source noise generated by the modulation error, etc.;  $\varepsilon_1$  is the channel noise due to the dithering of the relative phase and the depolarized scattering process of the local field, etc; *T* is the transmission efficiency of the channel.

In order to achieve a precise secret key rate, the parameters of the system are calibrated in detail. The losses at Alice's site do not affect the system performance because the signal level is set at Alice's output port. The transmission efficiency of the quantum channel (30.2 km single mode fiber) is 0.251. At Bob's site the transmission coefficient of the PBS is  $\eta_1 = 89.5\%$ . The detection efficiency of the pulse homodyne detector with the 50/50 coupler is  $\eta_2 = 66\%$ . Thus the total efficiency of Bob's setup is  $\eta = \eta_1\eta_2 = 59.1\%$ .

To find the optimum amplitude of the signal field,<sup>[10]</sup> we plot the secret key rate as a function of the signal field amplitude with different excess noises at the distance of 30 km, as shown in Fig. 3. It is clear that the optimum signal amplitude  $\alpha$  remains almost unchanged for different excess noises. For a fixed excess noise, the corresponding curve is flat near the optimum amplitude point. In our experiment, the signal amplitude was adjusted to be around 0.45.

Assuming that the attenuation of the single mode fiber is constant and linear, the amplitude of the signal field sent by Alice can be determined from the measured signal amplitude at Bob's site. At this stage, the experimental excess noise can be determined. With the reconciliation efficiency  $\beta = 80\%$  and the theoretical results in Refs. [10, 12], the security key rate per pulse can be given for each 80 blocks, as shown in Fig. 4. For each of the three points, the corresponding excess noise is less than 0.01. The estimated secret key rate for our system can be as much as 1 kbit/s at the current pulse repetition rate of 500 kHz. It is found that there were fluctuations for the measured values of excess noise and sometimes the excess noise can exceed 0.02. There are several factors which can contribute to the excess noises, including the source noise, modulation error, finite size effect, and the dithering of the relative phase between the signal and local fields,<sup>[11,20,26]</sup> etc. Besides the above-mentioned factors, we find that leakage from the local field into the signal path due to the depolarized scattering process of the local field can also contribute to the excess noise. In the following, we will concentrate on the issues of relative phase fluctuations and local field leakage.



Fig. 3. The secret key rate per pulse for different single field amplitudes and excess noises. The circle marks the maximum value of each curve, the asterisks are the experimental results,  $\alpha$  is the amplitude of the signal field.

The dithering of the relative phase is mainly due to the 80 m delay fiber. When we connect Alice's setup and Bob's setup directly without including the 80 m fiber and 30 km single mode fiber, the varying period of the relative phase between the signal and local fields is about 20–50 s. For a calibration period of 100 ms, we can lock the relative phase to be  $\pm 0.54^{\circ}$  (peak to peak value) using an active servo system. After inserting the 80 m fiber into the signal and local paths, the relative phase can be stabilized to be  $\pm 1.98^{\circ}$ . The reason is that the varying period of the relative phase is now reduced to 2–5 s while the calibration period is still 100 ms, so performance of the locking servo system will degrade. After inserting the 30 km single mode fiber, the relative phase can be locked to  $\pm 2.88^{\circ}$ . The contribution of the relative phase dithering to the excess noise can be estimated by a simple model.<sup>[20]</sup> Due to the phase variations, Bob's measurement basis will change from  $\hat{X}$  to  $\hat{X}'$  with  $\hat{X}' = \hat{X}\cos\theta_0 + \hat{P}\sin\theta_0$ ; here  $\theta_0$  is the phase dithering centered on zero. The added excess noise can be given by  $\langle (\Delta \hat{X}')^2 \rangle - \langle (\Delta \hat{X})^2 \rangle \approx \langle (\Delta \hat{P})^2 \rangle \theta_0^2 = V_A \theta_0^2$ ,

by employing the experimental value of  $\theta = 2.88\pi/180$ and  $V_A = 0.405$ ; the added excess noise is 0.001. As a result, due to the low signal modulation amplitude, the discrete modulation CVQKD protocol is quite robust against the phase locking noise. To improve the locking accuracy of the relative phase further, one can shorten the pulse width and increase the pulse repetition rate. By doing so the length of time-multiplexing fibers can be reduced and the varying period of the relative phase will increase accordingly, and at the same time the calibration period of the locking system can also be decreased effectively.

Due to the birefringence of the 30 km single mode fiber, the polarization state of the signal and local fields will drift slowly with the temperature.<sup>[27]</sup> Such drift can be solved by utilizing a dynamic polarization controller. Besides the low frequency drift, it is found that there also exist fast fluctuations of the polarization state of the local beam at frequency up to 40 kHz. Such fluctuations will lead to photon leakage from the local beam into the signal field and contribute to the excess noise.



Fig. 4. The measured spectrum of PHD output for different lengths of single mode fibers when the signal path at Alice's setup is disconnected. A: electronic noise, B: 30 km fiber (local power:  $1.89 \times 10^6$  photons per local pulse), C: 1 m fiber, D: 10 km fiber, E: 30 km fiber; the local power for C, D and E is  $5.7 \times 10^6$  photons per local pulse.

To analyze the above photon leakage effect, we disconnect the signal path at Alice's setup and analyze the spectrum of PHD output for different lengths of single mode fibers (1 m, 10 km and 30 km), as shown in Fig. 4. It is clear that the added excess noise is directly proportional to the fiber length and the power of the local field. For the 1-m-long fiber, the measured spectrum is fairly flat and it is identical to the spectrum of the vacuum field which is measured by disconnecting the signal path at Bob's setup. When longer single mode fibers  $(10 \,\mathrm{km} \text{ and } 30 \,\mathrm{km})$  are employed, a noise peak centered on zero frequency appears on the noise spectrum. This means that the signal field is not a vacuum field now and some depolarized local light is leaked into the signal path. Such leakage is most probably attributed to the depolarized Rayleigh scattering process.<sup>[28]</sup> When the local power is such that the PHD has a 10 dB shot noise to electronic noise ratio (the photons per local pulse is around  $5.7 \times 10^6$ ), the introduced excess noise can be above the level of 0.01. In our current setup, by decreasing the local beam's power to the level that the ratio of shot noise to electronic noise is around 6 dB (the photons per local pulse is around  $1.89 \times 10^6$ ), the excess noise can be reduced to the level of 0.005.

To realize a long-distance secure quantum transmission, the scattering process analyzed above must be taken into account and two approaches can be pursued to decrease the influence of such noise. Firstly, one can improve the extinction ratio of optical pulses. Secondly, a high shot noise to electronic noise ratio PHD with ultra-low electronic noise can be designed and utilized, in this case, the system can operate at a lower local power level and the unwanted leakage can be suppressed effectively.

In summary, a fiber based four-state discrete modulation continuous variable quantum key distribution system has been demonstrated. A secret key rate of 1 kbit/s is achieved at a transmission distance of 30.2 km. Various factors that lead to the excess noises are analyzed, with emphasis on the influence of the relative phase fluctuations and local field scattering process. Lastly, some approaches are proposed to reduce the excess noises due to the scattering process.

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010305-3

## Chinese Physics Letters

Volume 30 Number 1 January 2013

**GENERAL** 

- 010201 A Method of Choosing the Optimal Number of Singular Values in the Inverse Laplace Transform for the Two-Dimensional NMR Distribution Function JIANG Zhi-Min, WANG Wei-Min
- 010301 Double Barrier Resonant Tunneling in Spin-Orbit Coupled Bose–Einstein Condensates LI Zhi, WANG Jian-Zhong, FU Li-Bin
- 010302 An Alternative Approach to Construct the Initial Hamiltonian of the Adiabatic Quantum Computation

DUAN Qian-Heng, ZHANG Shuo, WU Wei, CHEN Ping-Xing

- 010303 One-Way Quantum Computation with Cluster State and Probabilistic Gate DIAO Da-Sheng
- 010304 Computation of Quantum Bound States on a Singly Punctured Two-Torus CHAN Kar-Tim, Hishamuddin Zainuddin, Saeid Molladavoudi
- 010305 Four-State Modulation Continuous Variable Quantum Key Distribution over a 30-km Fiber and Analysis of Excess Noise

WANG Xu-Yang, BAI Zeng-Liang, WANG Shao-Feng, LI Yong-Min, PENG Kun-Chi

- 010306 Fractals in Quantum Information Process BI Feng, LI Chuan-Feng
- 010501 Efficiency at Maximum Power of a Quantum Dot Heat Engine in an External Magnetic Field ZHANG Yan-Chao, HE Ji-Zhou
- 010502 Nonergodic Brownian Motion in a Collinear Particle-Coupled Harmonic Chain Model LU Hong, BAO Jing-Dong
- 010503 Single-Hopf Bursting in Periodic Perturbed Belousov–Zhabotinsky Reaction with Two Time Scales

LI Xiang-Hong, BI Qin-Sheng

- 010601 Accuracy Evaluation of NIM5 Cesium Fountain Clock LIU Nian-Feng, FANG Fang, CHEN Wei-Liang, LIN Ping-Wei, WANG Ping, LIU Kun, SUO Rui, LI Tian-Chu
- 010701 Research of Infrared Imaging at Atmospheric Pressure Using a Substrate-Free Focal Plane Arrav WU Jian-Xiong, CHENG Teng, ZHANG Qing-Chuan, ZHANG Yong, MAO Liang, GAO Jie,

CHEN Da-Peng, WU Xiao-Ping

### THE PHYSICS OF ELEMENTARY PARTICLES AND FIELDS

011201 Photoproduction of Large Transverse Momentum Dimuonium  $(\mu^+\mu^-)$  in Relativistic Heavy Ion Collisions YU Gong-Ming, LI Yun-De

O12101 Cluster Structure in Be Isotopes within Point-Coupling Covariant Density Functional TANG Zhong-Hua, LI Jia-Xing, JI Juan-Xia, ZHOU Tao
ATOMIC AND MOLECULAR PHYSICS
O13101 Variational-Integral Perturbation Corrections for Hydrogen Atom Content of the Atom Conte

- ZHAO Yun-Hui, PAN Yi-Qing, LI Wen-Juan, DENG Xia, HAI Wen-Hua
- Jong Zhang-Y 013201 Visible Light Emission in Highly Charged Kr<sup>17+</sup> Ions Colliding with an Al Surface YANG Zhi-Hu, XU Qiu-Mei, GUO Yi-Pan, WU Ye-Hong, SONG Zhang-Yong

013701 Trapping, Transporting, and Splitting Cold Molecules Employing a Spatial Liquid Crystal Modulator

```
GONG Tian-Lin, HUANG Yun-Xia, MU Ren-Wang, JI Xian-Ming, YANG Xiao-Hua
```

- 013702 Demonstration of Cold <sup>40</sup>Ca<sup>+</sup> Ions Confined in a Microscopic Surface-Electrode Ion Trap CHEN Liang, WAN Wei, XIE Yi, WU Hao-Yu, ZHOU Fei, FENG Mang
- 013703 Cooling and Crystallization of Trapped <sup>113</sup>Cd<sup>+</sup> Ions for Atomic Clock WANG Shi-Guang, ZHANG Jian-Wei, MIAO Kai, WANG Zheng-Bo, WANG Li-Jun

## FUNDAMENTAL AREAS OF PHENOMENOLOGY (INCLUDING **APPLICATIONS**)

- 014201 Birefringence Optical Feedback with a Folded Cavity in HeNe Laser WU Yun, TAN Yi-Dong
- 014202 Research on High-Intensity Picosecond Pump Laser in Short Pulse Optical Parametric Amplification PAN Xue, PENG Yu-Jie, WANG Jiang-Feng, LU Xing-Hua, OUYANG Xiao-Ping, CHEN Jia-Lin, JIANG You-En, FAN Wei, LI Xue-Chun
- 014203 Effect of Laser Pulse Width on the Laser Lift-off Process of GaN Films CHEN Ming, ZHANG Jiang-Yong, LV Xue-Qin, YING Lei-Ying, ZHANG Bao-Ping
- 014204 Spectral Characteristic Based on Fabry–Pérot Laser Diode with Two-Stage Optical Feedback WU Jian-Wei, Bikash NAKARMI
- 014205 Principal State Analysis for a Compact in-Line Fiber Polarization Controller LI Zheng-Yong, WU Chong-Qing, WANG Zhi-Hao, QIN Tao, WANG Yi-Xu
- 014206 Magnetic Field Induced Spectroscopy of <sup>88</sup>Sr Atoms Probed with a 10 Hz Linewidth Laser LIN Yi-Ge, WANG Qiang, LI Ye, LIN Bai-Ke, WANG Shao-Kai, MENG Fei, ZHAO Yang, CAO Jian-Ping, ZANG Er-Jun, LI Tian-Chu, FANG Zhan-Jun
- 014207 High Efficiency Grating Coupler for Coupling between Single-Mode Fiber and SOI Waveguides ZHANG Can, SUN Jing-Hua, XIAO Xi, SUN Wei-Min, ZHANG Xiao-Jun, CHU Tao, YU Jin-Zhong, YU Yu-De
- 014208 Two-Crystal Design and Numerical Simulations for High-Average-Power Second-Harmonic Generation

ZHONG Hai-Zhe, YUAN Peng, ZHU He-Yuan, QIAN Lie-Jia

- 014209 Stop Band Gap in Periodic Layers of Confined Atomic Vapor/Dielectric Medium LI Yuan-Yuan, LI Li, LU Yi-Xin, ZHANG Yan-Peng, XU Ke-Wei
- 014501 Effect of Interstitial Media on Segregation in Vertically Vibrated Granular Mixtures YUAN Xiao-Xian, LI Liang-Sheng, WEN Ping-Ping, SHI Qing-Fan, ZHENG Ning
- 014701 Dynamic Characteristics of Gas Transport in Nanoporous Media SONG Hong-Qing, YU Ming-Xu, ZHU Wei-Yao, ZHANG Yu, JIANG Shan-Xue

#### PHYSICS OF GASES, PLASMAS, AND ELECTRIC DISCHARGES

- 015201 Super-X Divertor Simulation for HCSB-DEMO Conception Design ZHENG Guo-Yao, PAN Yu-Dong, FENG Kai-Ming, HE Hong-Da, CUI Xue-Wu
- 015202 Influence of Discharge Voltage on Charged Particles in a Multi-Dipole Device in the Presence of an Ion Collecting Surface M. K. Mishra, A. Phukan, M. Chakraborty

### CONDENSED MATTER: STRUCTURE, MECHANICAL AND THERMAI PROPERTIES

us for Low The 016101 Formation of Co-implanted Silicon Ultra-Shallow Junctions for Low Thermal Budget Applications

Rehana Mustafa, S. Ahmed, E. U. Khan

016202 Temperature Effects of Electrorheological Fluids Based on One-Dimensional Calcium and **Titanium Precipitate** 

YAN Ren-Jie, WU Jing-Hua, LI Cong, XU Gao-Jie, ZHOU Lu-Wei

016203 Dynamic Mechanical Behavior and Failure Mechanism of Polymer Composites Embedded with Tetraneedle-Shaped ZnO Whiskers

RONG Ji-Li, WANG Dan, WANG Xi, LI Jian, XU Tian-Fu, LU Ming-Ming, CAO Mao-Sheng

016801 Electric-Field Switching of Bright and Dark Excitons in Semiconductor Crossed Nanowires LI Xiao-Jing, K. S. Chan

#### CONDENSED MATTER: ELECTRONIC STRUCTURE, ELECTRICAL, MAGNETIC, AND OPTICAL PROPERTIES

- 017101 Electronic Structure, Lattice Dynamics and Thermoelectric Properties of PbTe from **First-Principles Calculation** LU Peng-Xian, QU Ling-Bo
- 017102 Ultrafast and Broadband Terahertz Switching Based on Photo-Induced Phase Transition in Vanadium Dioxide Films CHEN Zhi, WEN Qi-Ye, DONG Kai, SUN Dan-Dan, QIU Dong-Hong, ZHANG Huai-Wu
- 017103 First-Principles Calculation of Lithium Adsorption and Diffusion on Silicene
- HUANG Juan, CHEN Hong-Jin, WU Mu-Sheng, LIU Gang, OUYANG Chu-Ying, XU Bo 017201 Spin-Dependent Electron Transport in an Armchair Graphene Nanoribbon Subject to Charge and Spin Biases

ZHANG Xiao-Wei, ZHAO Hua, SANG Tian, LIU Xiao-Chun, CAI Tuo

017202 Efficiency Enhancement of MEH-PPV:PCBM Solar Cells by Addition of Ditertutyl Peroxide as an Additive

LI Yan-Fang, YANG Li-Ying, QIN Wen-Jing, YIN Shou-Gen, ZHANG Feng-Ling

- 017301 On the Voltage and Frequency Distribution of Dielectric Properties and ac Electrical Conductivity in Al/SiO<sub>2</sub>/p-Si (MOS) Capacitors Ahmet Kaya, Semsettin Altındal, Yasemin Şafak Asar, Zekayi Sönmez
- 017302 High Deep-Ultraviolet Quantum Efficiency GaN P–I–N Photodetectors with Thin P-GaN **Contact Layer**

LIAN Hai-Feng, WANG Guo-Sheng, LU Hai, REN Fang-Fang, CHEN Dun-Jun, ZHANG Rong, ZHENG You-Dou

017303 Electronic Properties of a Phenylacetylene Molecular Junction with Dithiocarboxylate Anchoring Group

LIU Wen, XIA Cai-Juan, LIU De-Sheng

- 017401 Intra-Valley Spin-Triplet p + ip Superconducting Pairing in Lightly Doped Graphene ZHOU Jian-Hui, QIN Tao, SHI Jun-Ren
- 017402 Experimental Investigation of the Electronic Structure of Ca<sub>0.83</sub>La<sub>0.17</sub>Fe<sub>2</sub>As<sub>2</sub> HUANG Yao-Bo, RICHARD Pierre, WANG Ji-Hui, WANG Xiao-Ping, SHI Xun, XU Nan, WU Zheng, LI Ang, YIN Jia-Xin, QIAN Tian, LV Bing, CHU Ching-Wu, PAN Shu-Heng, SHI Ming, DING Hong
- 017501 Magnetoelastic Anisotropy of FeSiB Glass-Coated Amorphous Microwires LIU Kai-Huang, LU Zhi-Chao, LIU Tian-Cheng, LI De-Ren
- 017601 Detecting Larmor Precession of a Single Spin with a Spin-Polarized Tunneling Current GUO Xiao-Dong, DONG Li, GUO Yang, SHAN Xin-Yan, ZHAO Ji-Min, LU Xing-Hua
- 017801 Dosimetric Characteristics of a LKB:Cu,Mg Solid Thermoluminescence Detector Yasser Saleh Mustafa Alajerami, Suhairul Hashim, Ahmad Termizi Ramli, Muneer Aziz Saleh, Ahmad Bazlie Bin Abdul Kadir, Mohd. Iqbal Saripan
- 017802 Phase Shift of Polarized Light after Transmission through a Biaxial Anisotropic Thin Film HOU Yong-Qiang, LI Xu, HE Kai, QI Hong-Ji, YI Kui, SHAO Jian-Da
- a ranowire to a ranowire to a ranowire to a ranowire to a 017901 Current Density-Sensitive Welding of a Semiconductor Nanowire to a Metal Electrode TAN Yu, WANG Yan-Guo

017902 Elimination of the Schottky Barrier at an Au-ZnSe Nanowire Nanocontact via in Situ Joule Heating

TAN Yu, WANG Yan-Guo

# CROSS-DISCIPLINARY PHYSICS AND RELATED AREAS OF SCIENCE AND TECHNOLOGY

- 018101 Fabrication of Thin Graphene Layers on a Stacked 6H-SiC Surface in a Graphite Enclosure DENG Peng-Fei, LEI Tian-Min, LU Jin-Jun, LIU Fu-Yan, ZHANG Yu-Ming, GUO Hui, ZHANG Yi-Men, WANG Yue-Hu, TANG Xiao-Yan
- 018102 Phase Structure and Electrical Conduction of  $CaTi_{1-x}Sc_xO_{3-\delta}$  Ceramics ZHANG Qi-Long, LIU Yang, YANG Hui
- 018103 Characterization of Modified Tapioca Starch in Atmospheric Argon Plasma under Diverse Humidity by FTIR Spectroscopy P. Deeyai, M. Suphantharika, R. Wongsagonsup, S. Dangtip
- 018501 A SQUID Bootstrap Circuit with a Large Parameter Tolerance ZHANG Guo-Feng, ZHANG Yi, Hans-Joachim Krause, KONG Xiang-Yan, Andreas Offenhäusser, XIE Xiao-Ming
- 018701 Enhanced Response to Subthreshold Signals by Phase Noise in a Hodgkin–Huxley Neuron KANG Xiao-Sha, LIANG Xiao-Ming, LÜ Hua-Ping
- 018901 A Micro-Community Structure Merging Model Using a Community Sample Matrix LI Lin, PENG Hao, LU Song-Nian, TIAN Ying

#### GEOPHYSICS, ASTRONOMY, AND ASTROPHYSICS

- 019401 Large Bi-Polar Signature in a Perpendicular Electric Field of Two-Dimensional Electrostatic Solitary Waves Associated with Magnetic Reconnection: Statistics and Discussion LI Shi-You, ZHANG Shi-Feng, DENG Xiao-Hua, CAI Hong
- 019601 Spatial Distribution and Anisotropy of Energetic Particles Accelerated by Shock Waves: Focused Transport Model

ZUO Ping-Bing, FENG Xue-Shang

